# BV and BFV formalism beyond perturbation theory

Based on joint works w/ Benini-Safronov [2104.14886] and Benini-Pridham [2201.10225]

What is derived geometry?

- Traditional geometric frameworks, such as manifolds or schemes, are incapable to describe certain important geometric objects:
- (i) Quotients by non-free group actions:



and ignores in how many ways points get identified, e.g.

(ii) Non-transversal intersections:

$$X_{1} \underset{Y}{\times} X_{z} \text{ is in general singular, e.g.}$$

$$Y = \mathbb{R}^{3}, \quad X_{1} = \begin{cases} (x_{1}y_{1}z) \in \mathbb{R}^{3} : z = 0 \end{cases}, \quad X_{z} = \begin{cases} (x_{1}y_{1}z) \in \mathbb{R}^{3} : (xy)^{2} = z \end{cases}$$

$$\longrightarrow \quad X_{1} \underset{Y}{\times} X_{z} = \begin{cases} (x_{1}y_{1}z) \in \mathbb{R}^{3} : z = 0 \text{ and } (xy)^{2} = 0 \end{cases}$$

and ignores intersection multiplicities, e.g.







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Derived (algebraic) geometry resolves these issues by introducing a refined and powerful concept of space called derived stacks.

To get some intuition, we have to recall functors of points:

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schemes and more (what you get from glueing)

Interpretation of a functor

•  $\chi(\mathbb{R}^\circ) = "$ points in  $\chi''$ 

• 
$$X(\mathbb{R}^{1}) =$$
 "curres in  $X''$ 

•  $X(\mathbb{R}^2) =$  "surfaces in X''.



Derived stacks have a richer functor of points:



BRST/BV/BFV/... = perturbative/formal aspects of derived geometry:



Los-algebra (or DGC Aly)

Application 1: Nonperturbative (classical and finite dim.) BV formalism

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Physical scenario and motivation:

A physical system is typically described by

1.) a space of fields X

2.) an action of the gauge group  $X \times G \longrightarrow X$ 

3.) a gauge-invariant action function  $S: X \longrightarrow \mathbb{K}$ 

Want to determine the space of extrema of S modulo gauge symmetries, which involves taking quotients  $\chi/G$  and intersections  $\mathcal{L}^{dR}S = \mathcal{O}$ .

DAG potentially important for this problem!

### Mathematical formalization:

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Consider a smooth affine scheme  $X = S_{Pec}A$  with an action  $X \times G \longrightarrow X$  of a smooth affine group scheme  $G = S_{Pec}H$ .

A gauge-invariant function  $S: X \longrightarrow |K|$  is the same datum as a function S:  $[X/G] \longrightarrow |K|$  on the quotient stack

$$[X/G] := colim (X \rightleftharpoons XxG \rightleftharpoons XxG^2 \rightleftharpoons ...) \in dSt.$$

The derived critical locus is defined as the pullback in JSH

$$\frac{dCrit(S)}{dRS} \xrightarrow{V[X]}{V[X]} \begin{bmatrix} X/G \end{bmatrix} \qquad This solves the intersectionproblem  $d^{dR}S = 0$  and  
also counts multiplicities  
and stabilizers  $P$$$

Theorem (Benini-Safronov-AS):

 $dCrit(S) \simeq \left[\frac{2}{G}\right] \text{ is a derived quotient stack with } Z = \operatorname{Spec} O(Z).$ the derived affine scheme specified by the function dg-algebra  $O(Z) = \operatorname{Sym}_{A} \left( A \otimes \operatorname{gc-ZI}_{autifuelds} \oplus \operatorname{Tac-II}_{autifuelds} \right) \in \operatorname{dg}(Alg_{ZO})$  $\operatorname{drat}_{fields} \operatorname{drat}_{form} \operatorname{drat}_{form}$ 

#### Remarks:

- (1)  $\mathcal{JCri}(S)$  carries a canonical (-1)-shifted symplectic structure that can be computed via Lagrangian intersections.
- (2) The function dg-algebra can be compared with the BV formalism:

 $O(dC_{i}+(s)) \simeq O([Z/G]) \simeq \operatorname{Tot}^{\mathbb{T}} N^{\bullet}(G, O(Z))$   $Vot \simeq vou Est vou$ 

(3) The derived stack  $\mathcal{J}_{\mathcal{A}} + (S) \simeq [\mathcal{Z}/\mathcal{G}]$  is in general **not** affine, i.e. it is not determined by its function dg-algebra.

This is a new feature of the nonperturbative world!

A richer algebraic invariant is given by its dg-category of modules:

QCoh (dCrit (S)) ~ QCoh ([Z/G]) ~ O(Z), dy Mod G

Application 2: Nonperturbative BFV quantization

Physical scenario:

The phase space of 2nd-order gauge theories is a derived cotangent stack

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Similarly to before,  $p^{-1}(\sigma) = \operatorname{Spec} \mathcal{O}(p^{-1}(\sigma))$  is a derived affine scheme with

$$\mathcal{O}(\mu^{-1}(0)) = Sym_A(T_A + A \otimes g_{F-1}) \in dy(Aly_{\geq 0})$$

## Wanted:

Quantization of  $\mathcal{T}^*\mathcal{F}X/\mathcal{G}$  along the canonical 0-shifted Poisson structure. How does the resulting  $E_o = pointed$  dg-category look like:

$$\operatorname{QGh}(T^*[X/G])_{\mathrm{H}} = 2$$

# Strategy:

Turn Pridham's abstract deformation theoretic arguments into a concrete construction! Let me sketch the key ideas:

<u>Step 1:</u> Resolve T\*[×/6] ⊂ [r<sup>-1</sup>()/6] by a diagram of Lie algebra quotients:

This turns the global problem into a family of local stacky affine problems:

$$CE^{\bullet}(q, O(p^{-1}(\omega))) \longrightarrow CE^{\bullet}(q \oplus q, O(p^{-1}(\omega) \times G)) \xrightarrow{\rightarrow} \cdots$$

Step 2: Quantize level-wise via differential operators

$$CE^{\bullet}(q, \mathcal{O}(p^{-1}(\omega))) \xrightarrow{h} \longrightarrow CE^{\bullet}(q \oplus q, \mathcal{O}(p^{-1}(\omega) \times G)) \xrightarrow{h} \xrightarrow{\rightarrow} \cdots$$

and pass over to dg-categories of modules:

$$( \mathbf{F} := \left( (\mathbf{y}, \mathcal{O}(\mu^{-1}(\mathbf{o}))) \right)_{\mathbf{h}} d\mathbf{y} M d \xrightarrow{=} (\mathbf{y} \mathbf{e} \mathbf{y}, \mathcal{O}(\mu^{-1}(\mathbf{o}) \times \mathbf{G})) d\mathbf{y} M d \xrightarrow{=} (\mathbf{y} \mathbf{e} \mathbf{y}, \mathcal{O}(\mu^{-1}(\mathbf{o}) \times \mathbf{G})) d\mathbf{y} M d \xrightarrow{=} (\mathbf{y} \mathbf{e} \mathbf{y}, \mathcal{O}(\mu^{-1}(\mathbf{o}) \times \mathbf{G})) d\mathbf{y} M d \xrightarrow{=} (\mathbf{y} \mathbf{e} \mathbf{y}, \mathcal{O}(\mu^{-1}(\mathbf{o}) \times \mathbf{G})) d\mathbf{y} M d \xrightarrow{=} (\mathbf{y} \mathbf{e} \mathbf{y}, \mathcal{O}(\mu^{-1}(\mathbf{o}) \times \mathbf{G})) d\mathbf{y} M d \xrightarrow{=} (\mathbf{y} \mathbf{e} \mathbf{y}, \mathcal{O}(\mu^{-1}(\mathbf{o}) \times \mathbf{G})) d\mathbf{y} M d \xrightarrow{=} (\mathbf{y} \mathbf{e} \mathbf{y}, \mathcal{O}(\mu^{-1}(\mathbf{o}) \times \mathbf{G})) d\mathbf{y} M d \xrightarrow{=} (\mathbf{y} \mathbf{e} \mathbf{y}, \mathcal{O}(\mu^{-1}(\mathbf{o}) \times \mathbf{G})) d\mathbf{y} M d \xrightarrow{=} (\mathbf{y} \mathbf{e} \mathbf{y}, \mathcal{O}(\mu^{-1}(\mathbf{o}) \times \mathbf{G})) d\mathbf{y} M d \xrightarrow{=} (\mathbf{y} \mathbf{e} \mathbf{y}, \mathcal{O}(\mu^{-1}(\mathbf{o}) \times \mathbf{G})) d\mathbf{y} M d \xrightarrow{=} (\mathbf{y} \mathbf{e} \mathbf{y}, \mathcal{O}(\mu^{-1}(\mathbf{o}) \times \mathbf{G})) d\mathbf{y} M d \xrightarrow{=} (\mathbf{y} \mathbf{e} \mathbf{y}, \mathcal{O}(\mu^{-1}(\mathbf{o}) \times \mathbf{G})) d\mathbf{y} M d \xrightarrow{=} (\mathbf{y} \mathbf{e} \mathbf{y}, \mathbf{g} \mathbf{y}) d\mathbf{y} d \xrightarrow{=} (\mathbf{y} \mathbf{e} \mathbf{y}) d\mathbf{y} d \xrightarrow{=} (\mathbf{y} \mathbf{e} \mathbf{y}, \mathbf{g} \mathbf{y}) d\mathbf{y} d \xrightarrow{=} (\mathbf{y} \mathbf{e} \mathbf{y}) d \xrightarrow{=} (\mathbf{y} \mathbf{e} \mathbf{y}) d\mathbf{y} d$$

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autighosts)

<u>Step 3</u>: Obtain a global quantization by computing the homotopy limit (in dgCat) of the local quantizations:

#### Theorem (Benini-Pridham-AS):

For G reductive, the following dg-category is a model for  $QC_{L}(T^*TX/G1)_{L}$ . <u>Objects:</u>

Triples 
$$(\mathcal{E}_{0}, \nabla, \Psi)$$
 consisting of  
(1) a G-eqv.  $O(X)$ [[t]]-dg-module  $\mathcal{E}_{0}$  (wore functions)  
 $W/G$ -action)

(Z) a G-eqv. dy-connection 
$$\nabla: \mathcal{E}_{\bullet} \to \mathcal{N}^{1}(x)[\mathcal{I}_{\bullet}]] \otimes \mathcal{E}_{\bullet}$$
  
 $\Im(x)\mathcal{I}_{\bullet}\mathcal{I}_{\bullet}]$   
 $w.v.t. th d^{dR}$ , i.e.  $\nabla(\alpha s) = th d^{dR} \otimes s + \alpha \nabla(s)$   
 $\operatorname{CCR}: positions and momenta$ 

These data have to satisfy the following conditions:

(1) 
$$\nabla_{v} \nabla_{v} - \nabla_{v} \nabla_{v} = t_{v} \nabla_{v} \nabla_{v}$$
  
 $\nabla_{v} \Psi_{+} - \Psi_{+} \nabla_{v} = 0$   
 $\Psi_{+} \Psi_{+} + \Psi_{+} \nabla_{v} = 0$   
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Morphisms:

$$\frac{hom}{([\mathcal{E}_{\bullet},\nabla,\psi),(\mathcal{E}_{\bullet}',\nabla',\psi')]} := \frac{hom}{O(X)} \frac{G_{\bullet}\nabla_{\bullet}\psi}{O(X)} (\mathcal{E}_{\bullet},\mathcal{E}_{\bullet}')$$
preserving the  $G_{\bullet}\nabla_{\bullet}\psi$  structures strictly.

# Remark:

This result can be used to construct a dg-categorified lattice AQFT for non-Abelian Yang-Mills theory on directed graphs: